

X-ray induced electronic structure change in CuIr₂S₄

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The electronic structure of CuIr₂S₄ has been investigated using various bulk-sensitive x-ray spectroscopic methods near the Ir *L*₃-edge: resonant inelastic x-ray scattering (RIXS), x-ray absorption spectroscopy in the partial fluorescence yield (PFY-XAS) mode, and resonant x-ray emission spectroscopy (RXES). A strong RIXS signal (0.75 eV) resulting from a charge-density-wave gap opening is observed below the metal-insulator transition temperature of 230 K. The resultant modification of electronic structure is consistent with the density functional theory prediction. In the spin- and charge-dimer disordered phase induced by x-ray irradiation below 50 K, we find that a broad peak around 0.4 eV appears in the RIXS spectrum.

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The role of relativistic spin-orbit interactions in condensed matter systems is drawing much attention recently, and there are intensive research efforts to elucidate physics of materials with 5d electrons [1–5]. A thiospinel CuIr₂S₄ is such an example. The high temperature metallic phase has a cubic spinel structure [6], but the system becomes insulating and diamagnetic with triclinic structure below the metal-insulator transition (MIT) temperature of T_{MI} ∼ 226 K [7]. The “A”-site in the spinel structure is occupied by Cu, which is monovalent [8–10], and therefore inactive near the Fermi energy. The mixed-valent iridium ions (Ir^{3.5+}) occupy the B-sites, and form a pyrochlore lattice. From their powder neutron diffraction experiments, Radaelli and coworkers proposed an intriguing octamer model of charge ordering, in which Ir⁴⁺ ions form spin- and structural-dimers separated by non-magnetic Ir³⁺ ions [7]. This unusual charge-order (CO) was explained by Khomskii and Mizokawa as an orbitally driven Peierls transition. That is, a large overlap of Ir d orbitals along the (110)-type direction renders the band structure quasi-one-dimensional and susceptible to Peierls instability [11].

What is mystifying is its behavior when the sample is irradiated with x-rays or electrons. A significant drop in resistivity was observed when the sample was exposed to x-ray or electron irradiation at low temperatures (T < 50 K) [12, 13]. In addition, the resistivity change occurs quite slowly, over a period of several minutes, and this effect could be reversed by heating the sample above 100 K. Detailed single crystal x-ray scattering studies showed that the long-range CO was destroyed and a short-range incommensurate CO appeared in this irradiation-induced low-temperature (IILT) phase [14]. Although spectroscopic data suggest that substantial electronic structure modification occurs at the MIT, no spectroscopic changes in this IILT phase below 50 K have been reported to

date [15, 16]. For example, opening of an insulating gap of 0.15 eV below the MIT temperature was observed in both optics [18] and photoemission experiments [15], but the photoemission spectra did not show any change in the IILT phase [15, 16].

X-ray spectroscopy provides a natural avenue to investigate this intriguing behavior, since x-rays can act as a radiation source as well as a probing particle. However, detailed investigation of electronic structure using core level x-ray spectroscopy has been limited. Although several Ir 4f photoemission experiments have reported that Ir³⁺/Ir⁴⁺ charge disproportion does exist in the insulating phase, the high surface sensitivity of this technique obfuscates the interpretation [15–17]. While x-ray absorption spectroscopy (XAS) at the S *K*-edge has been valuable in detecting redistribution of S 3p states across the MIT [9], the Ir *L*₃-edge XAS has not been very useful due to broad spectral features arising from the short core-hole lifetime [9, 10]. In this work, we report bulk-sensitive x-ray spectroscopic studies which allow us to obtain direct information regarding the Ir 5d electronic states across the MIT and also in the IILT phase. Utilizing resonant inelastic x-ray scattering (RIXS) and resonant x-ray emission spectroscopy (RXES), we observe that the Ir 5d *t*_{2g} band shifts as a result of the opening of an insulating gap in the CDW phase, which confirms density functional theory (DFT) calculations [19–21]. We also find that a mid-infrared peak in the RIXS spectrum emerges around 0.4 eV in the IILT phase, suggesting that not only the crystal [14] but also the electronic structure are modified due to the x-ray irradiation. To explain this experimental observation, we propose a model in which local dimers freed from the long-range CO can hop around.

All x-ray measurements were carried out at the Advanced Photon Source on the undulator beamline 9ID-

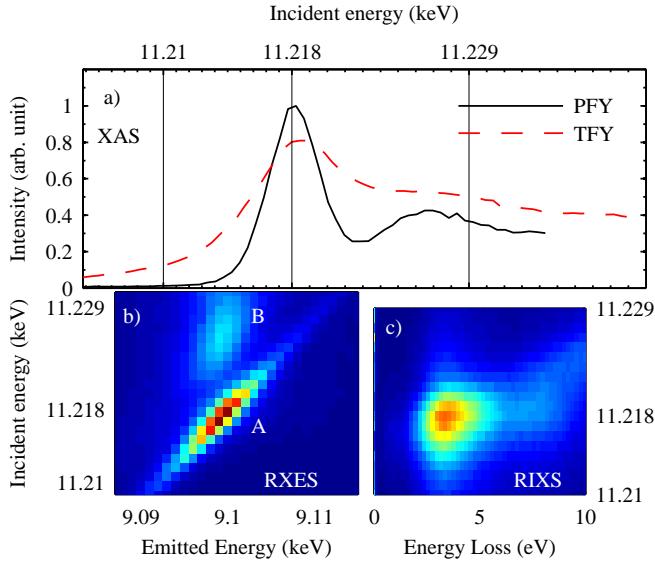


FIG. 1: (Color online) a) Ir L_3 -edge PFY-XAS taken by monitoring the intensity of La_2 ($3\text{d}_{3/2} \rightarrow 2\text{p}_{3/2}$) as a function of incident energy. Included is the normal XAS obtained by monitoring total fluorescence yield (TFY). b) and c) Wide range intensity map of both RXES and RIXS spectra as a function of the incident energy (E_i) and emitted energy (E_f) is shown. In Fig. 1(c), a RIXS intensity map is plotted as a function of E_i and energy loss ($\hbar\omega \equiv E_i - E_f$).

B. The beam was monochromatized by a double-crystal Si(111) and a Si(444) channel-cut secondary crystal. For the Ir L_3 -edge RIXS a spherical (1 m radius) diced Si(844) analyzer was used and an overall energy resolution of 0.15 eV (FWHM) was obtained. The incident photon polarization vector was rotated with a diamond phase plate to have it parallel to the vertical scattering plane (π polarization). In this geometry, due to the polarization factor, the contribution from Thomson scattering is minimized. At this photon energy, the sample probing depth is more than $10 \mu\text{m}$, allowing one to study true bulk electronic properties. Most of the measurements were carried out near $\mathbf{Q} = (7 7 8)$ in order to keep 2θ close to 90° . We also utilized a high-resolution setup with a Si(844) channel-cut secondary monochromator and a horizontal scattering geometry. In this setup no phase plate was needed to obtain π polarization. This provided an overall energy resolution of 0.08 eV (FWHM) and a significant reduction of the elastic line. For RXES and the XAS in the partial fluorescence yield mode (PFY-XAS) a spherical (1 m radius) diced Ge(337) analyzer was used to provide an overall energy resolution of 0.13 eV (FWHM). No diamond phase plate was used for these measurements. A single crystal of CuIr_2S_4 was grown by the bismuth solution method, as previously described in Ref. [14] and references therein. The crystal has a triangular shape with the surface normal along the (111) direction. Throughout this paper, we use the cubic notation ($a = 9.8474 \text{ \AA}$) for simplicity [6].

Wide energy range spectra obtained with all three experimental methods are shown in Fig. 1. In Fig. 1(a), the XAS spectrum near the Ir L_3 -edge is shown as a dashed line, which is obtained by monitoring total fluorescence yield. The PFY-XAS spectra shown as the solid line is the incident energy dependence of the La_2 emission line ($3\text{d}_{3/2} \rightarrow 2\text{p}_{3/2}$) intensity. The benefit of PFY mode is quite clear from this figure. By suppressing the spectral broadening due to 2p core-hole lifetime [22], sharper features are observed. The wide-range RXES and RIXS data are presented in the bottom panels of Fig. 1. In Fig. 1(b), the detailed evolution of the La_2 emission line as a function of incident energy (E_i) and emitted energy (E_f) is shown. In Fig. 1(c), a RIXS intensity map is plotted as a function of E_i and energy loss ($\hbar\omega \equiv E_i - E_f$).

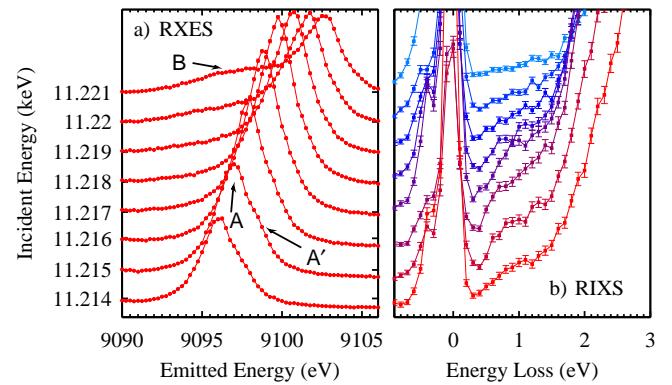


FIG. 2: (Color online) a) Incident energy dependence of RXES spectra taken at $T = 250$ K. b) Incident energy dependence RIXS spectra taken at $T = 220$ K and with $\mathbf{Q} = (6.3 7 8.3)$. Data was taken with incident photon energy of $11.214 \text{ keV} \leq E_i \leq 11.221 \text{ keV}$. The spectra have been shifted to match the scale on left.

The RXES data in Fig. 1(b) show two distinct features. There is a broad, strong feature (A) with maximum intensity around $E_i = 11.218$ keV and another feature (B) that resonates at higher incident energy. While feature B shows normal fluorescence behavior with fixed energy emission, the emitted photon energy of feature A shifts with E_i . This indicates that feature A comes from exciting an Ir 2p electron into an unoccupied *bound* state and thus the outgoing photon has a fixed energy loss. In this so-called Raman regime the RXES spectra can be understood as the unoccupied density of states, resonantly enhanced over the core hole width [23]. Feature A then corresponds to empty Ir 5d e_g states. The evolution of feature A as a function of E_i at $T = 250$ K is shown in Fig. 2(a). On the high energy side of A there is a prominent shoulder, A', which can be associated with empty t_{2g} states. Note that the A' feature is resonantly enhanced for E_i that is 2-3 eV below the resonance of feature A. The intensity ratio can be explained by the

fact that there are more empty e_g states than t_{2g} states (naively, 4 e_g states and 0.5 t_{2g} states). This intensity ratio, in addition to the 2p core-hole broadening, is why the t_{2g} state has not been revealed in previous XAS studies [9, 10] and in fact was not observed in our PFY-XAS. Only by utilizing the resonance enhancement of RXES can one see this state.

Fig. 3(a) shows the La_2 RXES spectra taken at $E_i = 11.215$ keV below and above the MIT. In the metallic phase, the La_2 emission line clearly has two components, A and A', while only A is seen in the insulating phase. The La_2 emission line at $T = 250$ K has been fitted with two Voigt functions, representing A and A'. The separation of the two peaks is estimated from the fit to be $\Delta E_{AA'} \approx 1.65$ eV [24]. As we go below the MIT ($T = 215$ K) the emission line appears as a single peak A; nevertheless we could extract A' using the same two-peak Voigt functions as for the $T = 250$ K case. Here we have assumed that the instrumental resolution and lifetime broadening are constant across the MIT. The result shows almost no shift for A but the A' feature shifts significantly by $\Delta E \approx 0.55$ eV. This observation can be consistently explained by the DFT calculation result, which is schematically shown in Fig. 3(b) [19, 21]. The main observation in our RXES study is that the opening of an insulating gap occurs due to a shift of the empty t_{2g} band with respect to the empty e_g band.

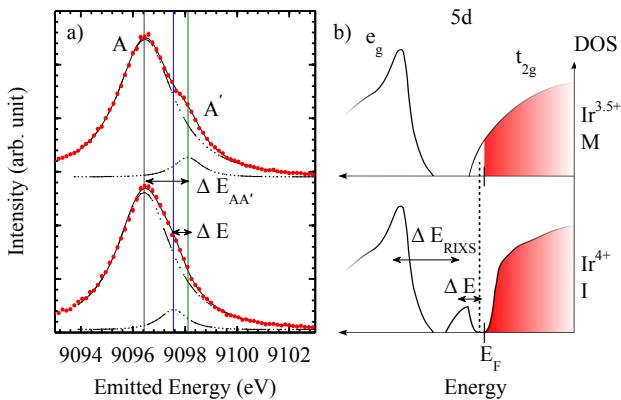


FIG. 3: (Color online) (a) La_2 RXES taken above and below the MIT at $E_i = 11.215$ keV. The dashed curves are the result of a fit (see text). ΔE represents the shift of A'. (b) Schematic diagram of the DOS below (I) and above (M) the MIT inferred from our experiment. ΔE_{RIXS} represents the splitting seen in RIXS.

In Fig. 1(c), a broad and strong RIXS feature in the 2-6 eV range can be seen, which shows large resonant enhancement near $E_i = 11.218$ keV. The weak higher energy (7-9 eV) feature is a fluorescence line. In between these high energy excitations and the elastic peak ($\hbar\omega = 0$) we observe a sharp edge-like RIXS feature around $\hbar\omega \approx 0.75$ eV, which is clearly seen in the individual scans shown in Fig. 2(b). This feature seems to resonate

at about 2 eV below the resonance energy of the 2-6 eV feature. This resonance behavior is crucial in identifying the observed spectral features. As in the interpretation of the RXES data, the intermediate states accessed through the main absorption feature at $E_i = 11.218$ keV are the Ir 5d e_g levels, and therefore the 2-6 eV excitation can be naturally associated with a transition into empty bands of hybridized Ir 5d e_g and S 3p character [8, 21]. The 0.75 eV feature, resonating around 11.216 keV, can then be identified as the d-d transition into an empty t_{2g} state. The energy difference between these two excitations represents the splitting between the empty e_g and t_{2g} states and is found to be about 2.8 eV. This splitting is also consistent with the DFT calculation value of about 2.6 eV [19].

In Fig. 4(a) we show in detail the temperature dependence of the low energy RIXS spectrum. The temperature dependence shows drastic spectral change as the sample goes through the MIT. In the metallic phase at $T = 250$ K, a non-zero intensity is observed at low energies. Once we go through the MIT a gap opens around 0.3 eV in the RIXS spectrum indicating the insulating state, which is consistent with the optical study [18]. According to the DFT calculation, in the high-temperature metallic phase, hybridized Ir t_{2g} and S 3p bands cross the Fermi level [19], giving rise to the low energy spectral weight observed below the gap. Below $T = 230$ K, an insulating gap opens in this band. The difference in the absolute level of the spectral weight above the gap is probably due to the difference in resonance condition, since electronic structure changes quite dramatically across the MIT. The momentum dependence of the insulating gap can be seen in Fig. 4(b). It was taken along high-symmetry directions as shown in the inset, where Γ is $\mathbf{Q} = (7 \ 7 \ 8)$. Scans at the high-symmetry positions, shown in Fig. 4(b), exhibit no \mathbf{q} dependence. This indicates that the gap opening is more or less uniform in reciprocal space, suggesting that the bandwidth of the empty t_{2g} band is fairly narrow. We have also measured the temperature and momentum dependence of the 2-6 eV feature (not shown), but no significant changes were seen.

When the sample was cooled below $T = 50$ K while irradiated with x-rays, a sudden drop of resistivity was observed [12, 13]. It was found that not only x-ray, but also electron, visible light, and even high energy ion beams could induce this IIIT phase [12, 14, 15, 25]. We have obtained high-resolution RIXS data in this phase, shown in Fig. 4(c), by cooling the system with x-rays impinging on the sample. A considerable increase of the spectral weight below 0.6-0.7 eV was observed. By subtracting the $T = 65$ K spectrum from the $T = 45$ K one, we can see that this intensity comes from the peak formed around 0.4 eV (shown in the inset) rather than from the tail of the elastic peak. The spectrum does not resemble the one at $T = 300$ K either, which indicates that the IIIT phase is different from the original high-temperature

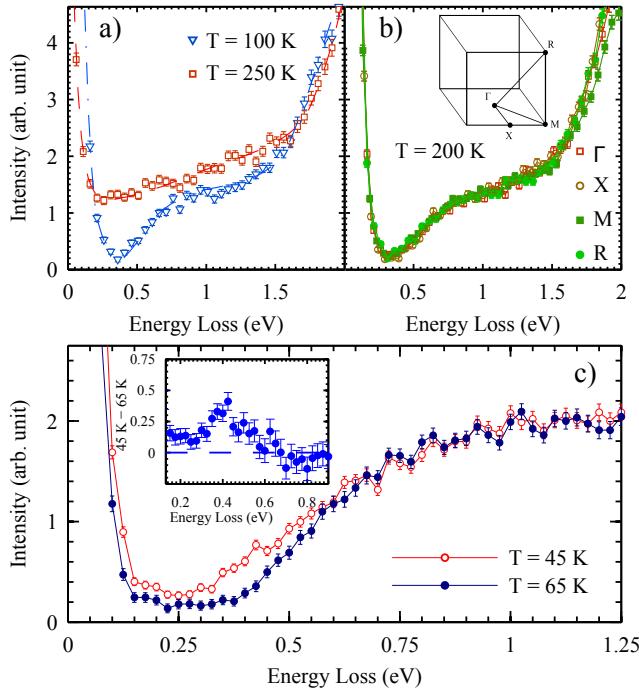


FIG. 4: (Color online) (a) Temperature dependence of the RIXS signal taken at $\mathbf{Q} = (6.15 \ 7 \ 8.5)$ for $E_i = 11.216$ keV. (b) Momentum dependence as measured at the high symmetry positions indicated in the inset. (c) High-resolution RIXS spectra taken above ($T = 65$ K) and within ($T = 45$ K) the charge disordered phase. To emphasize the difference induced by irradiation we show the subtracted spectra in inset.

metallic phase, in agreement with structural studies [26]. We also measured the La_2 RXES spectra for $T < 50$ K, but no difference was observed. We note that the larger elastic intensity ($\hbar\omega \lesssim 0.15$ eV) at $T = 45$ K compared to $T = 65$ K is due to the increased diffuse scattering when the long range CO is lost[14].

Even when the long-range CO is destroyed, it is known that individual dimers do exist in this phase [26]. One can imagine that when there is no long-range order, dimers can hop around more easily, and the increased hopping rate of dimers could be responsible for the decreased resistivity in this phase. In fact such dynamic dimers could be responsible for the mid-infrared (MIR) peak seen around 0.4 eV, which is reminiscent of the polaron peak observed in manganites [27]. Since dimers of Ir^{4+} - Ir^{4+} have significantly smaller bond lengths than Ir^{3+} - Ir^{3+} , they will introduce lattice distortion as well as charge-order defects to the crystal as they move around incoherently. Such an incoherent motion of dimers could be strongly coupled to the lattice, resulting in formation of a bipolaron-like object. Further high resolution spectroscopic study of this x-ray induced phase, such as optical conductivity, could shed light on this possibility.

In summary, we report comprehensive bulk-sensitive x-ray spectroscopy studies of electronic structure in

CuIr_2S_4 . We show that the Ir 5d t_{2g} band shifts as a result of the opening of a charge-density-wave gap, which is consistent with band theory calculations. In the low-temperature x-ray induced phase, we found that a broad peak around 0.4 eV appears in the RIXS spectrum, which could arise from the dynamic hopping of local dimers.

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- [1] Y. Okamoto, M. Nohara, H. Aruga-Katori and H. Takagi, Phys. Rev. Lett. **99**, 137207 (2007).
- [2] B.J. Kim *et al.*, Phys. Rev. Lett. **101**, 076402 (2008).
- [3] G. Jackeli and G. Khaliullin, Phys. Rev. Lett. **102**, 017205 (2009).
- [4] A. Shitade *et al.*, Phys. Rev. Lett. **102**, 256403 (2009).
- [5] D. Pesin and L. Balents, Nature Physics **6**, 376 (2010).
- [6] T. Furubayashi *et al.*, J. Phys. Soc. Jpn. **63**, 3333 (1994).
- [7] P. G. Radaelli *et al.*, Nature (London) **416**, 155 (2002).
- [8] J. Matsuno *et al.*, Phys. Rev. B **55**, R15979 (1997).
- [9] M. Croft *et al.*, Phys. Rev. B **67**, 201102 (2003).
- [10] N. Kijima *et al.*, J. Alloys Compd. **480**, 120 (2009).
- [11] D.I. Khomskii and T. Mizokawa, Phys. Rev. Lett. **94**, 156402 (2005).
- [12] H. Ishibashi *et al.*, Phys. Rev. B **66**, 144424 (2002).
- [13] T. Furubayashi *et al.*, Solid State Comm. **126**, 617 (2003).
- [14] V. Kiryukhin *et al.*, Phys. Rev. Lett. **97**, 225503 (2006).
- [15] K. Takubo *et al.*, Phys. Rev. Lett. **95**, 246401 (2005).
- [16] K. Takubo, T. Mizokawa, N. Matsumoto and S. Nagata, Phys. Rev. B **78**, 245117 (2008).
- [17] H.-J. Noh *et al.*, Phys. Rev. B **76**, 233106 (2007).
- [18] N.L. Wang *et al.*, Phys. Rev. B **69**, 153104 (2004).
- [19] T. Oda *et al.*, J. Phys. Condens Matter **7**, 4433 (1995).
- [20] T. Sasaki *et al.*, J. Phys. Soc. Jpn. **73**, 1875 (2004).
- [21] S. Sarkar, M. De Raychaudhury and T. Saha-Dasgupta, Phys. Rev. B **79**, 113104 (2009).
- [22] F.M.F. de Groot, M.H. Krisch and J. Vogel, Phys. Rev. B **66**, 195112 (2002).
- [23] M.H. Krisch *et al.*, Phys. Rev. Lett. **74**, 4931 (1995).
- [24] The value of $\Delta E_{AA'}$ is much smaller than the $e_g - t_{2g}$ splitting of 2.8 eV observed through RIXS. This difference is the RXES final state effect, since the 3d core hole interacts with the open 5d shell [S. Nakai *et al.*, J. Electron Spectrosc. Relat. Phenom. **137**, 363 (2004)]. However, the relative shift of A' with respect to A is not affected by the final state effect.
- [25] M. Koshimizu, H. Tsukahara and K. Asai, Nucl. Instr. and Meth. B **267**, 1125 (2009).
- [26] E.S. Božin, A.S. Masadeh, Y.S. Hor, J.F. Mitchell and S.J.L. Billinge, Phys. Rev. Lett. **106**, 045501 (2011).
- [27] K.H. Kim, J.H. Jung and T.W. Noh, Phys. Rev. Lett.

81, 1517 (1998).